## **Reminder of basic formula for symmetric catenary**

Catenary curve shape

$$y = c \cdot \cosh \frac{x}{c} \quad (m)$$
$$y' = \sinh \frac{x}{c}$$

Maximal sag

$$f_m = c \left( \cosh \frac{a}{2c} - 1 \right)$$

Wire length

$$l_s = 2c \sinh \frac{a}{2c}$$

Catenary constant

$$\frac{\sigma_{\rm H}}{\gamma} = c$$

## 3. Stress in wire

Basic facts:

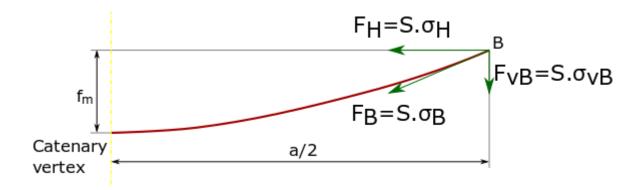
- Horizontal component of stress is constant along a catenary.

$$\sigma_H = const$$

- Vertical component of stress in a point at catenary is equal to the weight of weir between this point and catenary vertex.

$$\sigma_{\rm vB} = \gamma \cdot \frac{l_s}{2}$$

- Stress in any point at catenary has always tangential direction to the catenary.



Force in wire can be expressed by several ways:

1) By vector sum of its vertical and horizontal component:

$$F_B = S \cdot \sqrt{\sigma_H^2 + \sigma_{\rm vB}^2}$$

Using previous formulas for  $\sigma_{vB}$  and consequently for  $l_s$  and for catenary constant c, it can be rewritten to:

$$F_B = S \cdot \sigma_H \cdot \sqrt{1 + \left(\sinh\frac{a}{2c}\right)^2}$$

Which is equal to

$$F_B = S \cdot \sigma_H \cdot \cosh \frac{a}{2c}$$

2) The  $\cosh \frac{a}{2c}$  in the previous formula can be exchanged for  $y_B/c$  (see the formula for catenary curve shape and take into account that x is  $x_B=a/2$  for the point B)

$$F_B = S \cdot \sigma_H \cdot \frac{y_B}{c}$$

Using formula for catenary constant c:

 $F_B = S \cdot \gamma \cdot y_B$ 

It means that the force in point B is the same as weight of wire of length  $y_B$ 

3) The  $\cosh \frac{a}{2c}$  in the formula above for  $F_B$  can be exchanged for  $\frac{f_m}{c} + 1$  (See equation for maximal sag). Using formula for catenary constant c, following will be obtained:

$$F_B = S \cdot (\gamma \cdot f_m + \sigma_H)$$

It means, that the force in boint B is the same as sum of horizontal force and the weight of wire of the length of vertical distance between this point and the vertex (maximal sag).

## **4.** Equation of state

An equation for horizontal stresses under different temperature and overloading will be derived.

- Change of wire's length can due to a change of themperature reads

$$\Delta l_{\vartheta} = \alpha l_0 (\vartheta_1 - \vartheta_0)$$

where

index 0 – initial state (known)

index 1 – new state (computed)

- $\alpha$  coefficient of thermal expansion( °C<sup>-1</sup>),  $l_0$  initial wire's length (m),
- $\mathcal{G}_0$  initial wire's themperature (°C<sup>-1</sup>),
- $\mathcal{G}_1$  new wire's temperature (°C<sup>-1</sup>).
- Change of wire's length can due to a change of weight (overloading due to icing)

$$\Delta l_{\sigma} = \frac{l_0}{E} (\sigma_{H1} - \sigma_{H0})$$

where

E - Young's modulus (MPa),

σ<sub>H0</sub> - horizontal stress component during initial state (MPa),

 $\sigma_{\rm H1}$  - horizontal stress component during new state (MPa).

Overall change of wire's length as a sum of contributions:

$$\Delta l = l_1 - l_0 = \Delta l_{\mathcal{G}} + \Delta l_{\sigma} = l_0 \left[ \alpha \left( \mathcal{G}_1 - \mathcal{G}_0 \right) + \frac{1}{E} \left( \sigma_{\mathrm{H}1} - \sigma_{\mathrm{H}0} \right) \right]$$

Overall change of wire's length will now be expressed also from the catenary curve equation. Length of wire reads

$$l_s = 2c \sinh \frac{a}{2c}$$

Taking only first two elements of Thaylor series of this formula will give (it is also result of integration of parabolic approximation of catenary):

$$l_s = a + \frac{a^3 \gamma^2}{24 \sigma_H^2}$$
$$l_k = a + \frac{a^3 \gamma_k^2}{24 \sigma_{Hk}^2}$$

where

a - span (m).

 $\gamma$  - specific weight per 1 m of wire (MPa · m<sup>-1</sup>).

Therefore the overall change of wire's length is

$$\Delta l = l_1 - l_0 = \frac{a^3}{24} \left( \frac{\gamma_1^2}{\sigma_{\rm H1}^2} - \frac{\gamma_0^2}{\sigma_{\rm H0}^2} \right)$$

Using the two equations for the overall change of wire's length will give the equation of state

$$l_0 \left[ \alpha \left( \vartheta_1 - \vartheta_0 \right) + \frac{1}{E} \left( \sigma_{H1} - \sigma_{H0} \right) \right] = \frac{a^3}{24} \left( \frac{\gamma_1^2}{\sigma_{H1}^2} - \frac{\gamma_0^2}{\sigma_{H0}^2} \right)$$

We can usually consider the approximation  $l_0 = a$ 

and write

$$\alpha(\vartheta_1 - \vartheta_0) + \frac{1}{E}(\sigma_{H1} - \sigma_{H0}) = \frac{a^2}{24} \left(\frac{\gamma_1^2}{\sigma_{H1}^2} - \frac{\gamma_0^2}{\sigma_{H0}^2}\right)$$

After rearrangement that gives a cubic equation for the uknown  $\sigma_{H1}$ :

$$\sigma_{\rm H1}^{3} + \sigma_{\rm H1}^{2} \left[ \frac{E a^{2} \gamma_{0}^{2}}{24 \sigma_{\rm H0}^{2}} + \alpha E \left( \vartheta_{\rm I} - \vartheta_{\rm 0} \right) - \sigma_{\rm H0} \right] - \frac{a^{2} \gamma_{\rm I}^{2} E}{24} = 0$$

It is also common to express the specific weights  $\gamma_0 a \gamma_1$  using the specific weight of a pure conductor  $\gamma_v$  and an its overloading

$$\gamma_1 = \gamma_v \, z_1 \qquad \gamma_2 = \gamma_v \, z_2$$

The state equation than reads:

$$\sigma_{\rm H1}^{3} + \sigma_{\rm H1}^{2} \left[ \frac{E \gamma_{\rm v}^{2}}{24} \left( \frac{a z_{0}}{\sigma_{\rm H0}} \right)^{2} + \alpha E \left( \vartheta_{\rm I} - \vartheta_{\rm 0} \right) - \sigma_{\rm H0} \right] - \frac{E \gamma_{\rm v}^{2}}{24} \left( a z_{\rm I} \right)^{2} = 0$$